

# Poor Man's Advanced Analysis Techniques

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New Horizons in Lattice Field Theory

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# Advanced Fitting Methods

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“The fit of data by a sum of exponentials and in particular by real exponents may be very badly conditioned”

Lanczos 1956

Wiscombe & Evans 1977

Ruhe & Wedin 1980

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I will not offer proofs of any of the methods but rather introduce them practically for you to implement the proofs are left as an exercise for you

# Advanced Fitting Methods

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Designed for a-symmetric matrices of correlation functions  
(poor man = a-symmetric)

- Variational Projection
- Input-parameter free multi-exponential fitting
- Generalized Pencil of Function (GPOF)
- Matrix Prony

Very simple idea:  
if a fit function has mixed linear/non-linear dependence  
on the fit parameters, one does not need to perform a  
numerical minimization on all the parameters - one  
can first perform a linear least squares on the linear  
parameters, solving as a function of the non-linear  
ones

Very simple idea: first perform linear-least squares

in lattice QFT, correlation functions fit with

$$\chi^2 = \sum_{t,t'} [y(t) - f(\lambda, t)] C_{t,t'}^{-1} [y(t') - f(\lambda, t')]$$

$$f(\lambda, t) = \sum_n Z_n \lambda_n^t \quad \lambda_n = e^{-E_n} \quad \lambda = \{Z_n, \lambda_n\}$$

$$\frac{\partial \chi^2}{\partial Z_n} = 0$$

(repeated indices summed over)

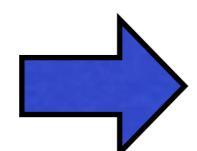
$$= \left[ -\lambda_n^t C_{t,t'}^{-1} (y(t') - Z_m \lambda_m^{t'}) - (y(t) - Z_m \lambda_m^t) C_{t,t'}^{-1} \lambda_n^{t'} \right]$$

Very simple idea: first perform linear-least squares

symmetry in  $t \leftrightarrow t'$

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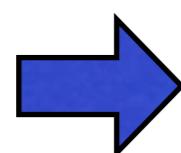
  $Z_m \lambda_m^t C_{t,t'}^{-1} \lambda_n^{t'} = y(t) C_{t,t'}^{-1} \lambda_n^{t'}$

Very simple idea: first perform linear-least squares

symmetry in  $t \leftrightarrow t'$

$$\frac{\partial \chi^2}{\partial Z_n} = 0$$

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$$Z_m \lambda_m^t C_{t,t'}^{-1} \lambda_n^{t'} = y(t) C_{t,t'}^{-1} \lambda_n^{t'}$$

$$Z_m \Lambda_{m,n} = y(t) C_{t,t'}^{-1} \lambda_n^{t'} \quad \Lambda_{m,n} = \lambda_m^t C_{t,t'}^{-1} \lambda_n^{t'}$$

Very simple idea: first perform linear-least squares

$$\frac{\partial \chi^2}{\partial Z_n} = 0 \quad \rightarrow \quad Z_m = \Lambda_{m,n}^{-1} y(t) C_{t,t'}^{-1} \lambda_n^{t'}$$

we have solved for the overlap factors as functions of the eigenvalues

plug these solutions back into chi^2 and perform numerical minimization on just the eigenvalues

Very simple idea: first perform linear-least squares

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When counting the degrees of freedom in the fit DON'T forget to count the overlap factors you have also determined!

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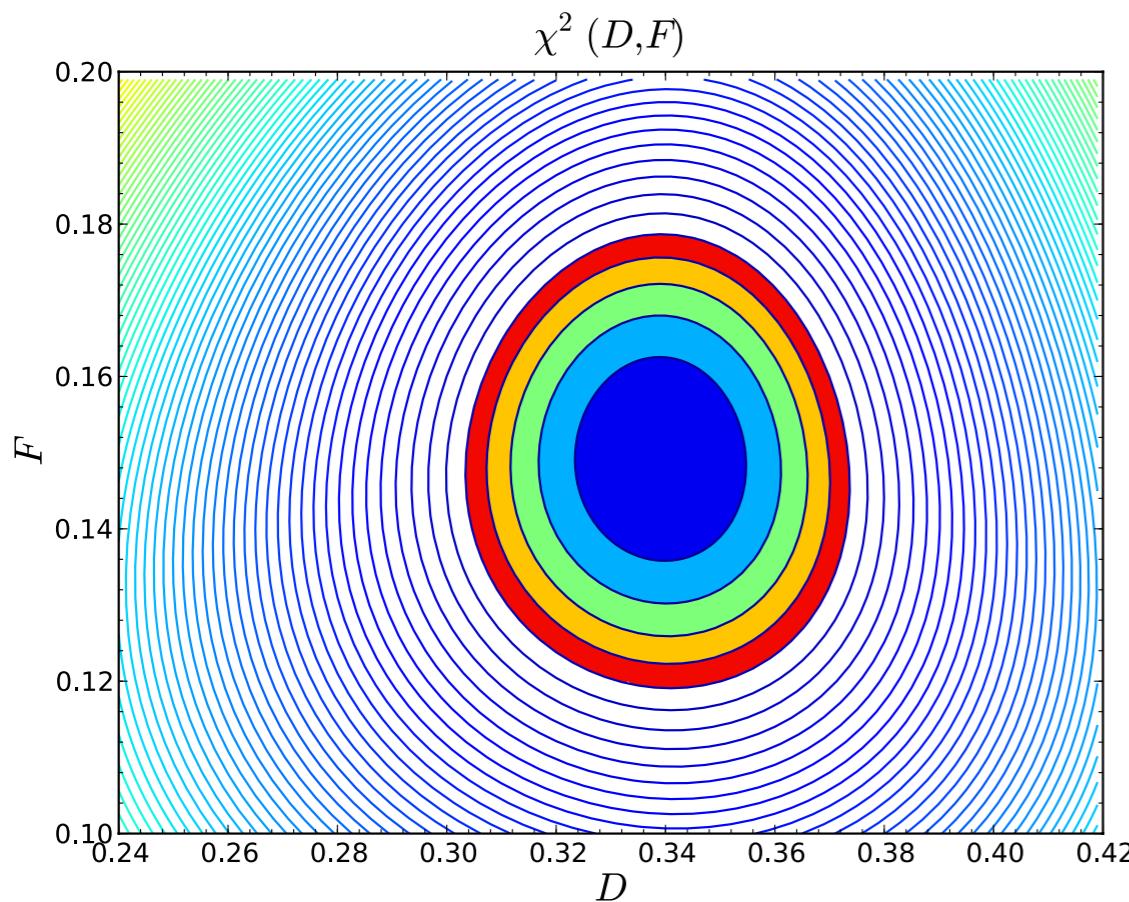
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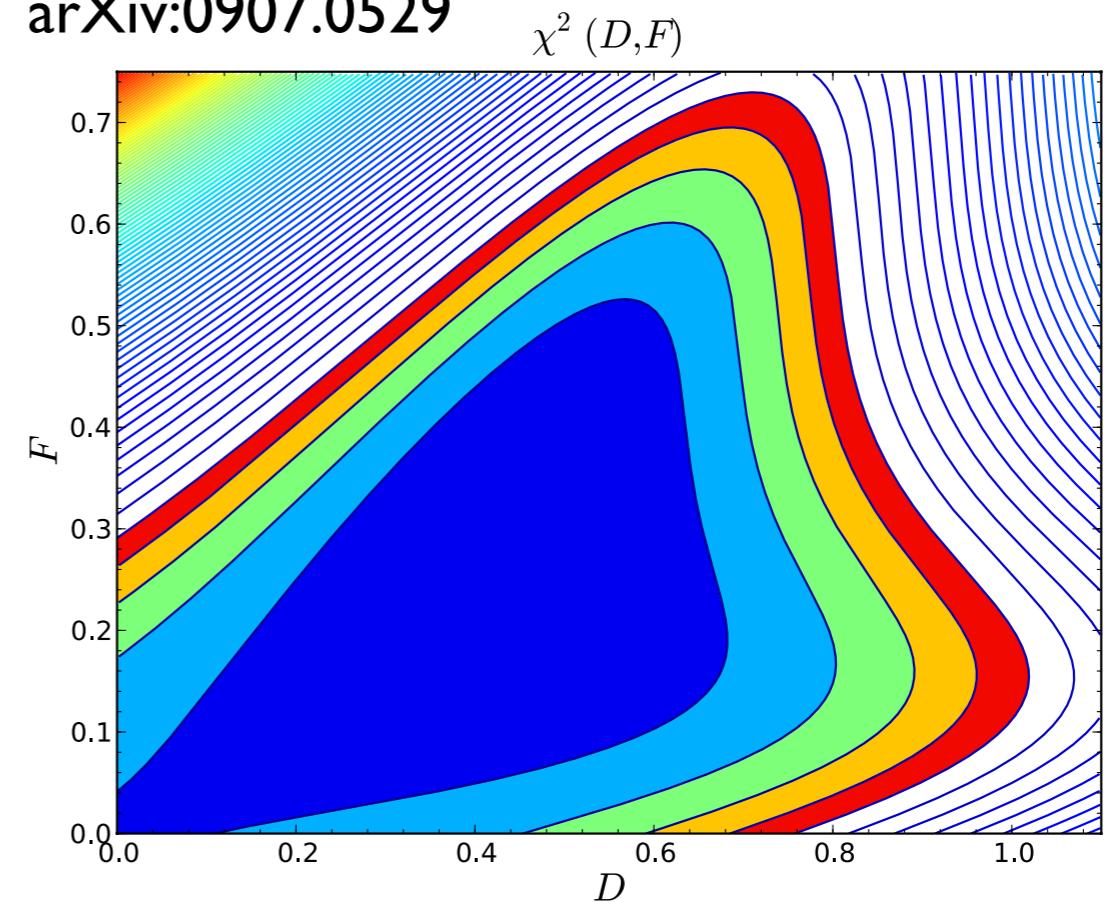
readily extend this to matrix of correlation functions

can also apply the same idea to chiral extrapolation formulae which usually have mixed linear/non-linear dependence on LECs

E. Jenkins, A. Manohar, J. Negele and AWL  
arXiv:0907.0529



NLO: 11 LECs



NNLO: 30 LECs

Fit to Octet and Decuplet baryon mass results with SU(3)  
Baryon Chiral Perturbation Theory

“The fit of a sum of exponentials to noisy data”  
(Note: typos in paper)

Given the outputs of a non-degenerate n-dimensional linear system

$$y_t = c^T x_t \quad c_i \neq 0, i = 1, \dots, n \quad \{x_t, c\} \in \mathbb{R}^n$$

$$x_{t+1} = T x_t \quad x_0^T = (1, \dots, 1)$$

$$y_t = c^T T^t x_0$$

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construct the  $p \times q$  Hankel Matrix  
with  $p, q > n$

$$H_{ij} = y_{i+j-1}$$

$$H = \begin{pmatrix} y_0 & \cdots & y_{q-1} \\ \vdots & \ddots & \vdots \\ y_{p-1} & \cdots & y_{q+p-2} \end{pmatrix}$$

The Hankel Matrix can be factorized

$$H = \begin{pmatrix} y_0 & \cdots & y_{q-1} \\ \vdots & \ddots & \vdots \\ y_{p-1} & \cdots & y_{q+p-2} \end{pmatrix} \quad y_t = c^T T^t x_0$$

$$H = FG = \begin{pmatrix} c^T \\ c^T T \\ \vdots \\ c^T T^{p-1} \end{pmatrix} (x_0 \ T x_0 \ \cdots \ T^{q-1} x_0)$$

The factorization matrices can be inverted  $H = FG$

$$F : \mathbb{R}^n \rightarrow \mathbb{R}^p \quad G : \mathbb{R}^q \rightarrow \mathbb{R}^n \quad (H \text{ is } p \times q \text{ matrix})$$

$$\mathbf{1}^n = \tilde{F}^{-1} H \tilde{G}^{-1}$$

$$\{\tilde{F}^{-1}, \tilde{G}^{-1}\} = \text{pseudo-inverse}\{F, G\}$$

The shifted Hankel Matrix can be factorized

$$H_{ij} = y_{i+j-1}$$

$$H_{ij}^+ = y_{i+j}$$

$$y_t = c^T T^t x_0$$

$$H = FG$$

$$H^+ = FTG$$

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We can construct a matrix similar to  $T$

$$T_s = \tilde{F}^{-1} H^+ \tilde{G}^{-1} = \tilde{F}^{-1} F T G \tilde{G}^{-1}$$

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How do we determine  $\{F, G\}$  ?

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- First - pick a window in time of interest you want to analyze,  $t_i, t_f$
- One can (should) shift with  $\Delta t = \tau > 1$

$$H_{ij}^\tau = y_{i+j-1+\tau}$$

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$$H = \begin{pmatrix} y_0 & \cdots & y_{q-1} \\ \vdots & \ddots & \vdots \\ y_{p-1} & \cdots & y_{q+p-2} \end{pmatrix}$$

```

1 # written in python language
2 import numpy as np
3
4 corr = # data array in form array([ncfg,nt,data[ncfg,nt]])
5 ncfg = corr.shape[0]
6 nt = corr.shape[1]
7 dt = 4 #shift you chose
8
9 h_tmp = np.zeros([nt/2 -dt, nt/2 -dt])
10 h_tmp_shift = np.zeros_like(h_tmp)
11 y_t = np.mean(corr, axis=0)
12 for t1 in range(nt/2 - dt):
13     for t2 in range(nt/2 - dt):
14         h_tmp[t1,t2] = y_t[t1+t2]
15         h_tmp_shift[t1,t2] = y_t[t1+t2+dt]
16
17 hankel = h_tmp[t0:tf:dt,0:tf-t0:dt]
18 hankel_shift = h_tmp_shift[t0:tf:dt,0:tf-t0:dt]
```

# Input Parameter Free Multi-Exp Fits

De Groen & De Moor  
Comp.Appl. Math. 20 (1987)

Perform a singular value decomposition on  $H$

$$H = U\Sigma V^T \quad \Sigma = \text{diag}(\text{singular values})$$

$$T_s = \Sigma^{-1/2} U^T H^+ V \Sigma^{-1/2} \quad H^+ = FTG$$

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One very important point: need to truncate the singular values to a reasonable number

- If you desire just the ground state - chose range of time suitably (contamination from just a single state) and take two singular values
- For determining 2 states (gs + 1st excited) pick 3 singular values etc.
- study the singular values to see how much information you can get from correlation function

# Input Parameter Free Multi-Exp Fits

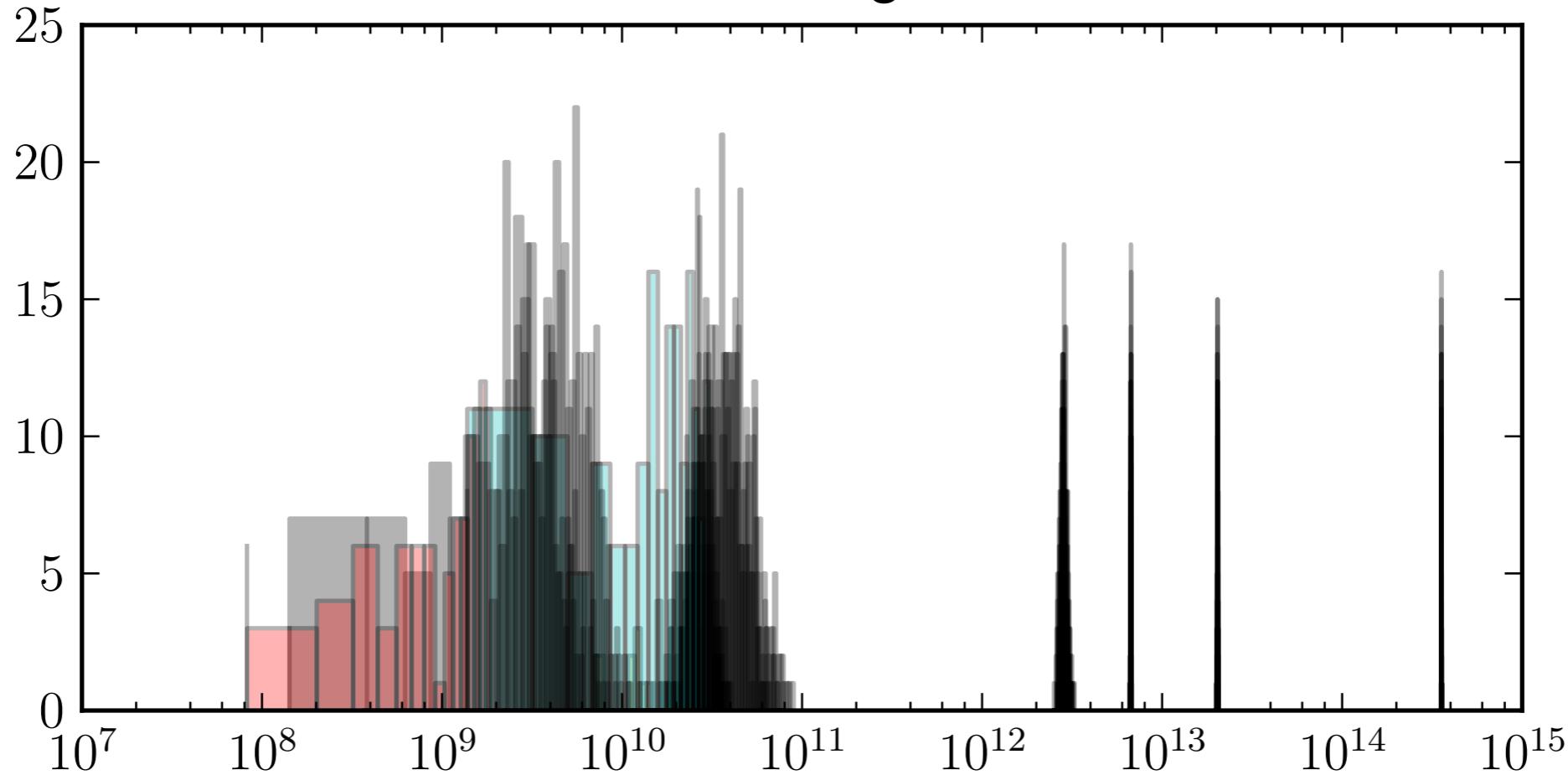
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Don't get greedy (don't try and determine too many excited states)

distribution of singular values



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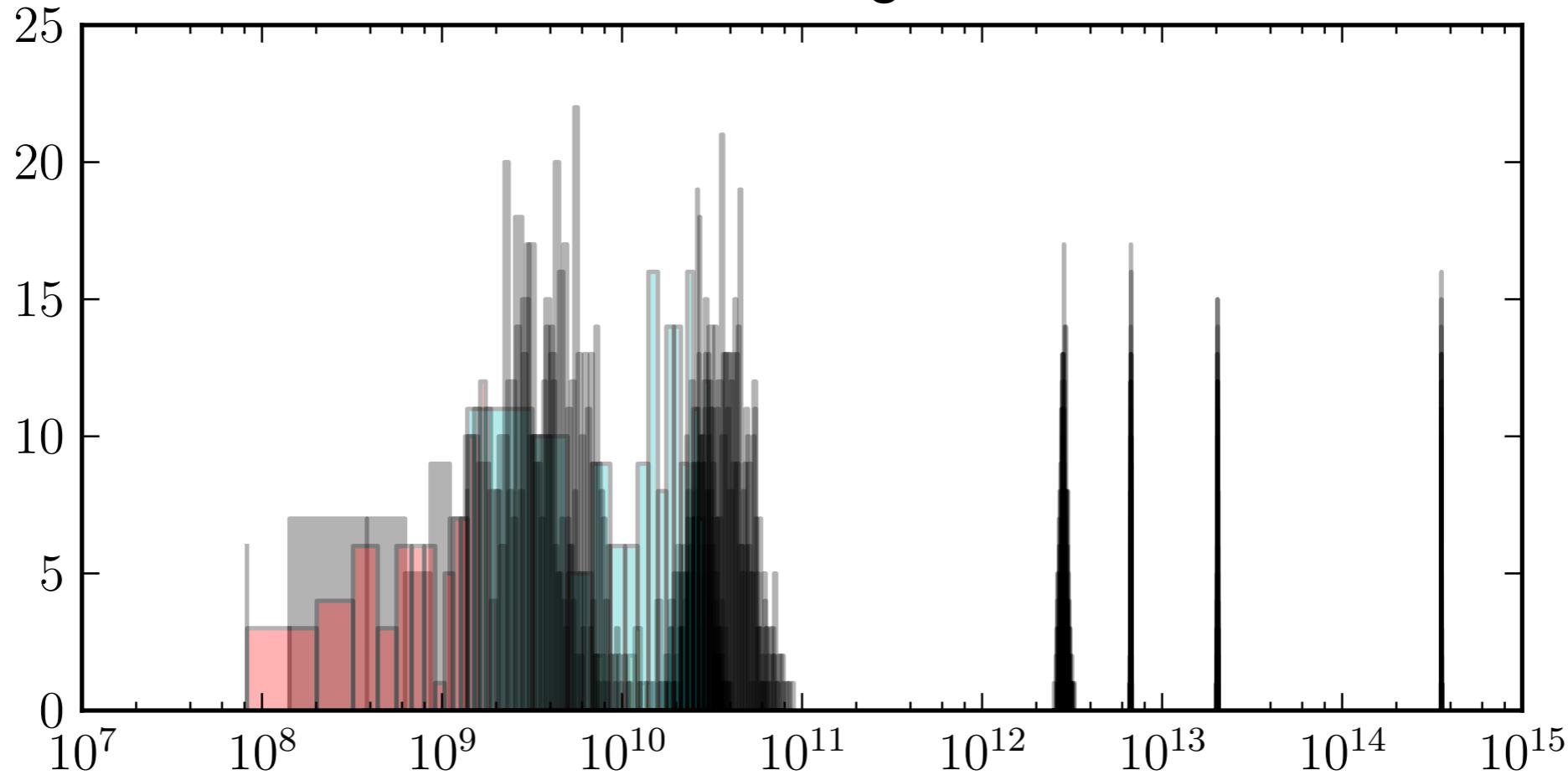
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in this case, at most, one can determine 4 energies and only 3 with confidence (the 4th absorbs the “slop”)

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take  $n_s$  largest singular values

$$T_s = \Sigma^{-1/2}[0 : n_s] U^T[0 : n_s, :] H^+ V[:, 0 : n_s] \Sigma^{-1/2}[0 : n_s]$$

$$H^+ = FTG$$

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$$T_s = \Sigma^{-1/2}[0 : n_s] U^T[0 : n_s, :] H^+ V[:, 0 : n_s] \Sigma^{-1/2}[0 : n_s]$$

$$\lambda_n, O_n = \text{eig}(T_s) \quad H^+ = FTG$$

$$E_n = -\ln(\lambda_n)$$

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Now we have the eigen-energies - how do we get the overlap factors?

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See VarPro!

# Input Parameter Free Multi-Exp Fits

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Comp.Appl. Math. 20 (1987)

$$H_{ij} = y_{i+j-1} \quad H_{ij}^+ = y_{i+j} \quad H = U\Sigma V^T \quad n_s$$

$$T_s = \Sigma^{-1/2} [0 : n_s] U^T [0 : n_s, :] H^+ V[:, 0 : n_s] \Sigma^{-1/2} [0 : n_s]$$

$$\lambda_n, O_n = \text{eig}(T_s) \quad E_n = -\ln(\lambda_n)$$

$$Z_m = \Lambda_{m,n}^{-1} y(t) C_{t,t'}^{-1} \lambda_n^{t'} \quad \Lambda_{m,n} = \lambda_m^t C_{t,t'}^{-1} \lambda_n^{t'}$$

Use these values to seed the multi-exponential fit!

This is a generalization of the previous idea (Hankel Matrix) extended to a vector or matrix of correlators.

For the sake of time - I refer you to the above references

Prony = some guys name

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G. R. de Prony Journal de l'cole Polytechnique, volume I, cahier 22, 24-76 (1795)

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Matrix = matrix

Start with a vector of correlation functions  
perhaps a single src and multiple sinks

$$y(t) = \begin{pmatrix} y_1(t) \\ y_2(t) \\ \vdots \\ y_N(t) \end{pmatrix}$$

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perhaps a single src and multiple sinks

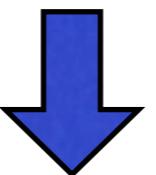
We would like to construct an  
operator such that

$$y(t) = \begin{pmatrix} y_1(t) \\ y_2(t) \\ \vdots \\ y_N(t) \end{pmatrix}$$

$$y(t + \tau) = \hat{T}(\tau)y(t)$$

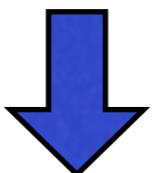
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$$y(t + \tau)y^T(t) = \hat{T}(\tau)y(t)y^T(t)$$

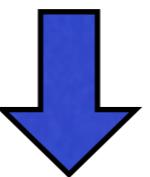
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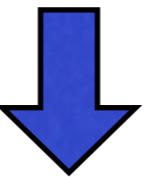
$$y(t + \tau)y^T(t) = \hat{T}(\tau)y(t)y^T(t)$$

Transpose

$$y(t + \tau) = \hat{T}(\tau)y(t)$$



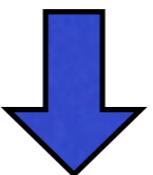
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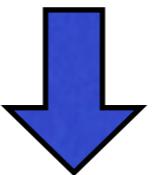
$$\hat{M}(\tau)y(t + \tau)y^T(t) = \hat{V}(\tau)y(t)y^T(t) \quad \hat{T} = \hat{M}^{-1}\hat{V}$$

Transpose

$$y(t + \tau) = \hat{T}(\tau)y(t)$$



$$y(t + \tau)y^T(t) = \hat{T}(\tau)y(t)y^T(t)$$

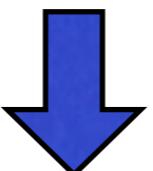


$$\hat{M}(\tau)y(t + \tau)y^T(t) = \hat{V}(\tau)y(t)y^T(t) \quad \hat{T} = \hat{M}^{-1}\hat{V}$$

Transpose

Nothing special about  $t$

$$y(t + \tau) = \hat{T}(\tau)y(t)$$



$$\hat{M}(\tau) \sum_{t_0}^{t_0 + \Delta t} y(t + \tau)y^T(t) = \hat{V}(\tau) \sum_{t_0}^{t_0 + \Delta t} y(t)y^T(t)$$

We have assumed T (and M and V) are independent of  $t$   
but this is precisely what we desire from a “transfer matrix”

$$\hat{M}(\tau) \sum_{t_0}^{t_0 + \Delta t} y(t + \tau) y^T(t) = \hat{V}(\tau) \sum_{t_0}^{t_0 + \Delta t} y(t) y^T(t)$$

a solution

$$\hat{M}(\tau) = \left[ \sum_{t_0}^{t_0 + \Delta t} y(t + \tau) y^T(t) \right]^{-1}, \quad \hat{V}(\tau) = \hat{V} = \left[ \sum_{t_0}^{t_0 + \Delta t} y(t) y^T(t) \right]^{-1}$$

must sum over sufficient number of time slices to make matrices “full rank” for two-components, must sum over at least two time slices

$$\hat{M}(\tau) \sum_{t_0}^{t_0 + \Delta t} y(t + \tau) y^T(t) = \hat{V}(\tau) \sum_{t_0}^{t_0 + \Delta t} y(t) y^T(t)$$

a solution

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most robust results come from maximizing  $\Delta t$   
 this requires that over a large range of time our ansatz  
 is satisfied - only  $N$  states contribute in  $t_0 \rightarrow t_0 + \Delta t$

$$y(t + \tau) = \hat{M}^{-1}(\tau) \hat{V} y(t) = \hat{T}(\tau) y(t)$$

$$\hat{M}(\tau) = \left[ \sum_{t_0}^{t_0 + \Delta t} y(t + \tau) y^T(t) \right]^{-1}, \quad \hat{V}(\tau) = \hat{V} = \left[ \sum_{t_0}^{t_0 + \Delta t} y(t) y^T(t) \right]^{-1}$$

given  $\mathbf{M}$  and  $\mathbf{V}$  one then solves the eigenvalue equation

$$\hat{T}(\tau) q_n = (\lambda_n)^\tau q_n \quad \lambda_n = e^{-E_n} \quad \hat{T}(\tau) = \hat{M}^{-1}(\tau) \hat{V}$$

given  $\lambda_n$  check to see if ansatz is satisfied - over range of  $t_0 \rightarrow t_0 + \Delta t$  there should be no significant evidence of excited state contamination

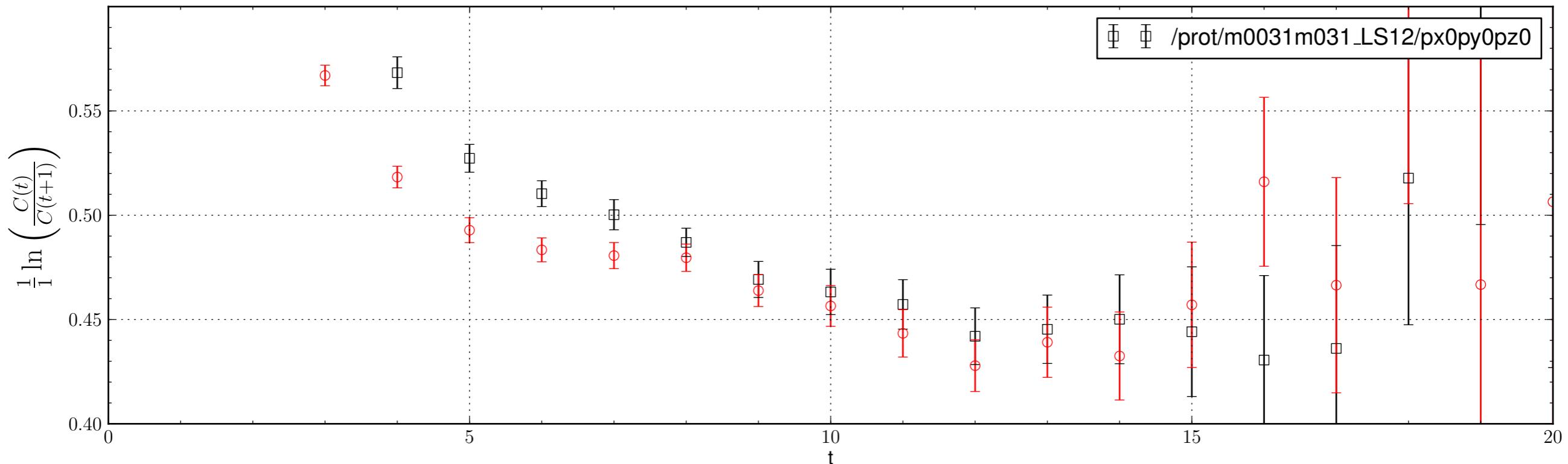
## in Python

```
20  ''' MATRIX PRONY '''
21  import numpy as np
22  from scipy import linalg as spla
23  def matrix_prony(corr,ti,tau,td=1):
24      corr = np.array(corr)
25      n_cfg,n_t,n_corr = corr.shape
26      if tau < n_corr:
27          print 'Matrix Prony Error:'
28          print ' must chose at least %d time slices' %n_corr
29          sys.exit(-1)
30      y_avg = np.mean(corr, axis=0)
31      Minvt = np.array(map(np.outer,np.roll(y_avg,-td,0),y_avg))
32      Vinvt = np.array(map(np.outer,y_avg,y_avg))
33      Minv = np.sum(Minvt[ti:ti+tau],axis=0)
34      Vinv = np.sum(Vinvt[ti:ti+tau],axis=0)
35      Transfer = np.dot(Minv,spla.inv(Vinv))
36      le,0e = spla.eig(Transfer)
37      ind = np.argsort(le); indr = ind[::-1]
38
39      return le, 0e, indr
```

$$y(t + \tau) = \hat{T}(\tau)y(t)$$

**take correlation function with  
one source and two sinks**

$$y(t) = \begin{pmatrix} y_{SP}(t) \\ y_{SS}(t) \end{pmatrix}$$

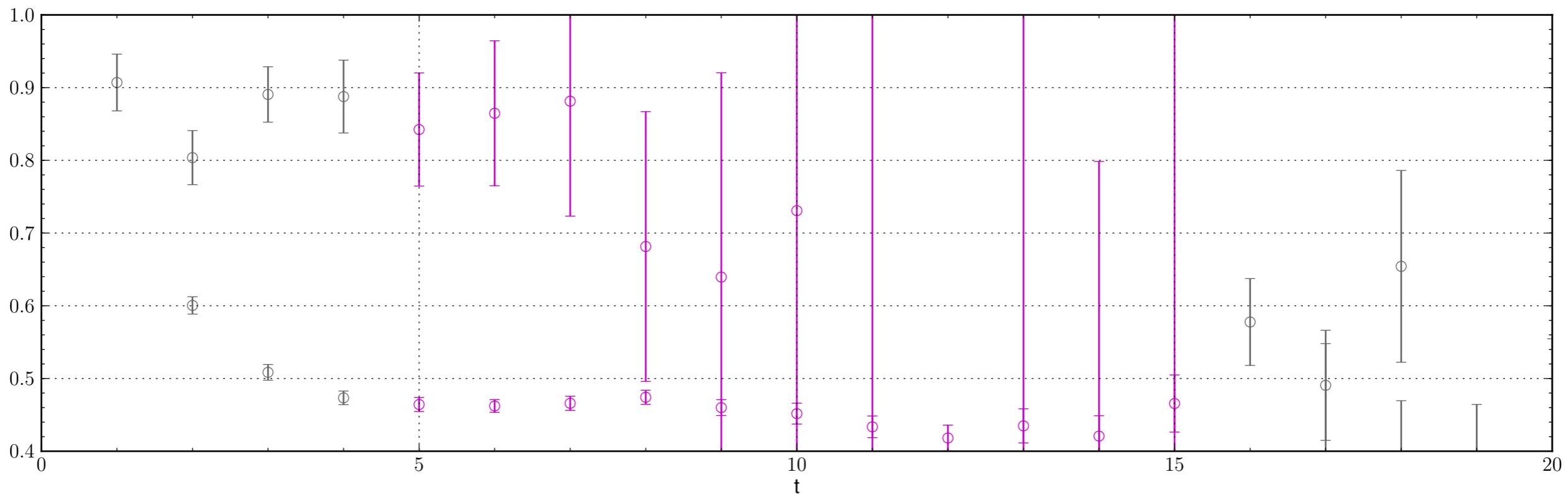


**between  $t=5$  and  $t=15$  there appear to be two states**

$$y(t + \tau) = \hat{T}(\tau)y(t)$$

**take correlation function with  
one source and two sinks**

$$y(t) = \begin{pmatrix} y_{SP}(t) \\ y_{SS}(t) \end{pmatrix}$$

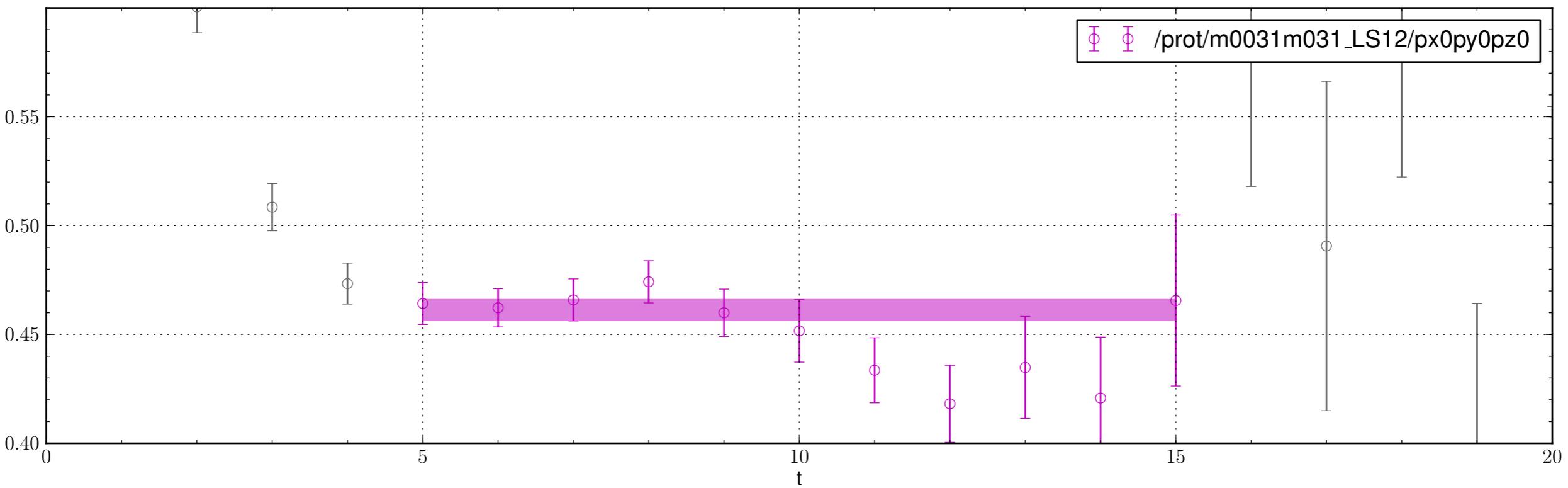


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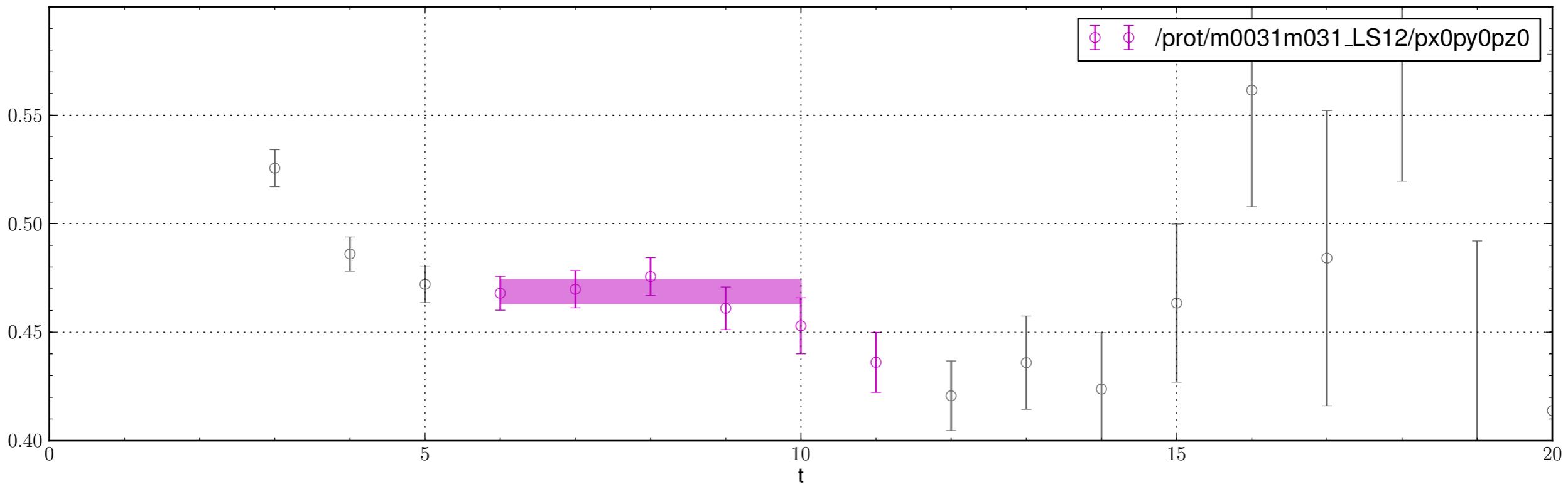


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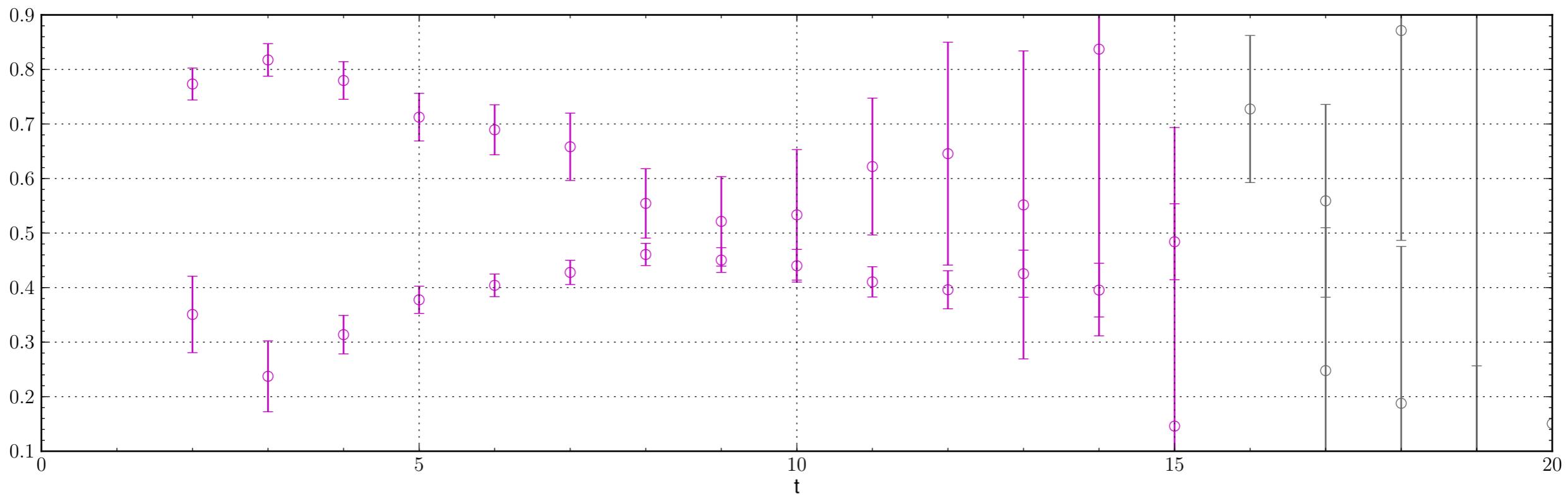
**between  $t=6$  and  $t=11$**

**also good**

$$y(t + \tau) = \hat{T}(\tau)y(t)$$

**take correlation function with  
one source and two sinks**

$$y(t) = \begin{pmatrix} y_{SP}(t) \\ y_{SS}(t) \end{pmatrix}$$



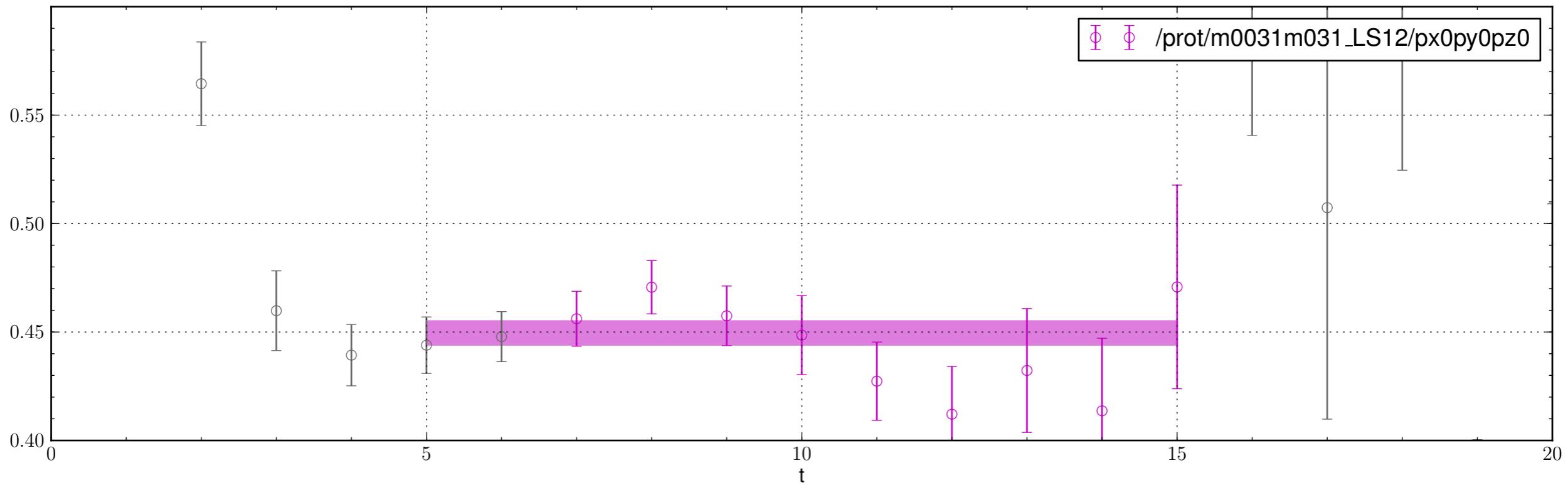
**between t=2 and t=15**

**too aggressive**

$$y(t + \tau) = \hat{T}(\tau)y(t)$$

**take correlation function with  
one source and two sinks**

$$y(t) = \begin{pmatrix} y_{SP}(t) \\ y_{SS}(t) \end{pmatrix}$$



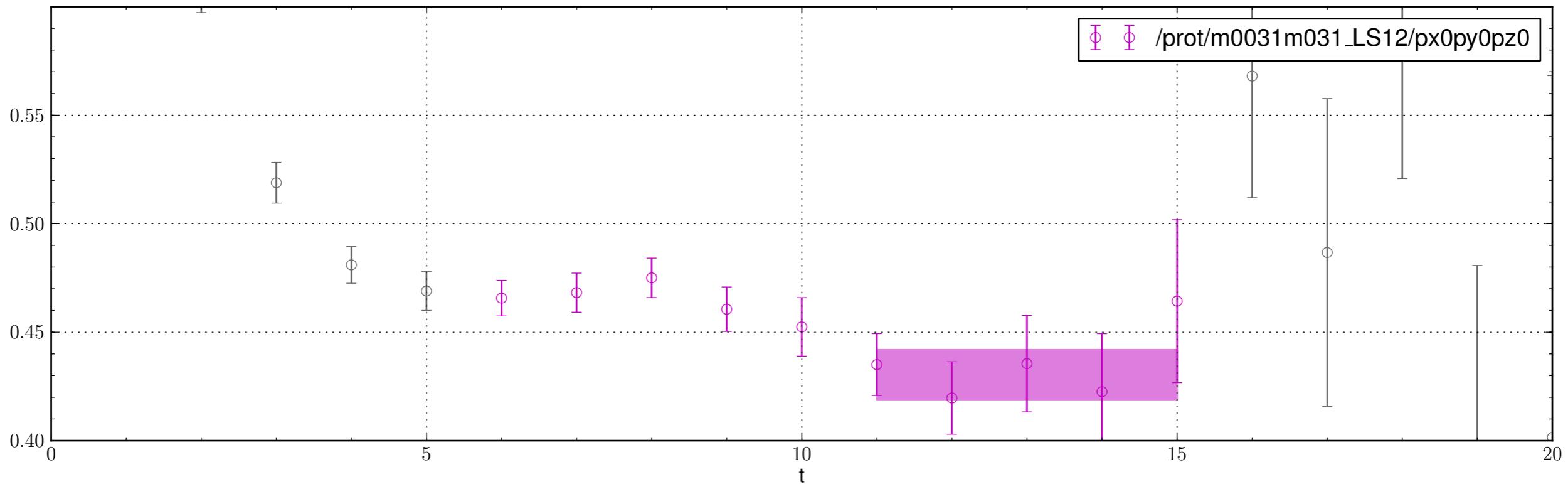
**between  $t=7$  and  $t=15$**

**also good**

$$y(t + \tau) = \hat{T}(\tau)y(t)$$

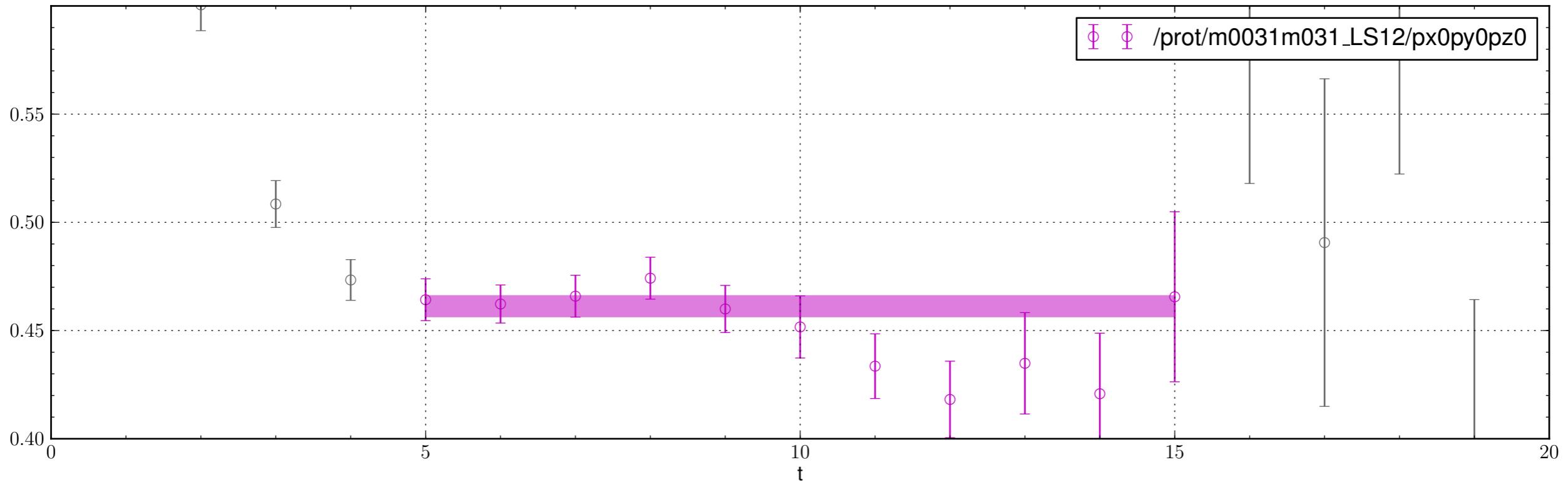
**take correlation function with  
one source and two sinks**

$$y(t) = \begin{pmatrix} y_{SP}(t) \\ y_{SS}(t) \end{pmatrix}$$



**between  $t=7$  and  $t=15$**

**different range of fit ok**



**exponential decay of signal-to-noise for baryons challenging**  
 is late-time dip of effective mass true ground state value?  
 lack of positive-definite correlation functions allow “false plateaus”  
 results much more sensitive to choices of fit-range, MP-range  
 want algorithm to weight all possible reasonable choices

## The algorithm I currently use (systematics are not uniquely defined)

- pick minimum  $\Delta t_{MP}$
- pick minimum  $\Delta t_{plat}$  in exp fit  $t_f = t_i + \Delta t_{plat}$   
 $\Delta t_{plat}$  choice guided by excited state masses
- pick minimum and maximum t you want to consider (correlator dependent)
- loop over independently chosen  $t_0, \Delta t_{MP}, t_i, \Delta t_{plat}$

$$\hat{M}(\tau) = \left[ \sum_{t_0}^{t_0 + \Delta t} y(t + \tau) y^T(t) \right]^{-1}$$

The algorithm I currently use  
(systematics are not uniquely defined)

- specific for baryons, for each fit, pick a weight factor

$$w_i = \frac{Q_i}{\sigma_i} \quad Q = \int_{\chi^2_{min}}^{\infty} d\chi^2 \mathcal{P}(\chi^2, dof) \quad Q \in [0, 1]$$

$\sigma_i$  = statistical uncertainty for a given fit

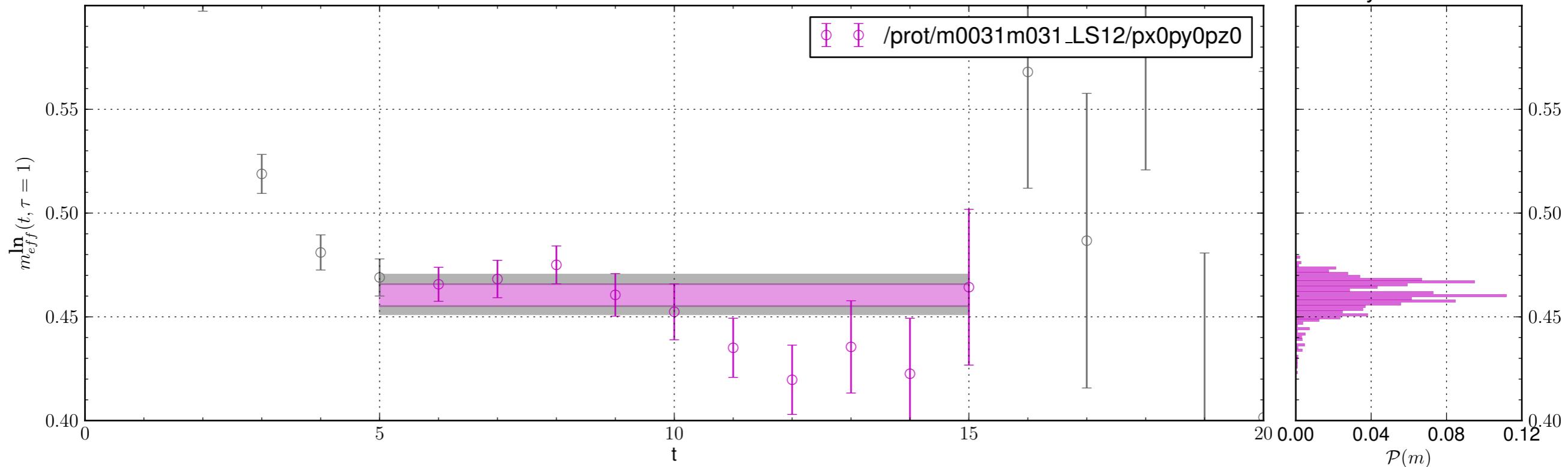
this choice allows for late time fluctuations being real, but suppresses them by their larger statistical uncertainty

$$\bar{m} = \frac{\sum_i w_i m_i}{\sum_j w_j}$$

$$y(t + \tau) = \hat{T}(\tau)y(t)$$

**take correlation function with  
one source and two sinks**

$$y(t) = \begin{pmatrix} y_{SP}(t) \\ y_{SS}(t) \end{pmatrix}_{\text{systematic}}$$



**16% and 84% quantiles chosen for systematic uncertainty  
inner band statistical outer band stat+sys added in quadrature**

$\mathcal{P}(m)$  = mass probability distribution function

$$\mathcal{P}(m_i) = \frac{w_i}{\sum_j w_j}$$

# Poor Man's Advanced Fitting Methods

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**Go forth and try these methods**