## Problem set 4 - Many-body theory SFI7534

## 1. Gaussian integral of complex variables.

(a) Let  $\mathbb{H}$  be an Hermitean positive definite  $n \times n$  matrix. Then show that

$$\int \prod_{i=1}^{n} \frac{\mathrm{d}(\Re z_i) \,\mathrm{d}(\Im z_i)}{\pi} e^{-(\mathbf{z}^*)^T \cdot \mathbb{H} \cdot \mathbf{z} + (\mathbf{z}^*)^T \cdot \mathbf{J} + (\mathbf{J}^*)^T \cdot \mathbf{z}} = (\mathrm{Det}\,\mathbb{H})^{-1} e^{(\mathbf{J}^*)^T \cdot \mathbb{H}^{-1} \cdot \mathbf{J}}$$

where  $\mathbf{J} \in \mathbb{C}^n$  is a complex vector.

(b) Use this result to show that

$$\langle z_{i1}^* \dots z_{im}^* z_{j1} \dots z_{jm} \rangle = \sum_{\text{all pairings}} \mathbb{H}_{j1,iP_1}^{-1} \dots \mathbb{H}_{jm,iP_m}^{-1},$$

where  $\{P_1, \ldots, P_m\}$  is a permutation of  $\{1, \ldots, m\}$ , and the average is defined as

$$\langle A \rangle = \frac{\int \prod_{i=1}^{n} \frac{\mathbf{d}(\Re z_{i})\mathbf{d}(\Im z_{i})}{\pi} e^{-(\mathbf{z}^{*})^{T} \cdot \mathbb{H} \cdot \mathbf{z}} A}{\int \prod_{i=1}^{n} \frac{\mathbf{d}(\Re z_{i})\mathbf{d}(\Im z_{i})}{\pi} e^{-(\mathbf{z}^{*})^{T} \cdot \mathbb{H} \cdot \mathbf{z}}}$$

## 2. Gaussian integral of Grassmann variables.

(a) Let  $\mathbb{H}$  be an Hermitean positive definite  $n \times n$  matrix (actually, this is not necessary). Then show that

$$\int \prod_{i=1}^{n} \mathrm{d}\bar{\eta}_{i} \mathrm{d}\eta_{i} e^{-(\bar{\boldsymbol{\eta}})^{T} \cdot \mathbb{H} \cdot \boldsymbol{\eta} + (\bar{\boldsymbol{\eta}})^{T} \cdot \boldsymbol{\xi} + \left(\bar{\boldsymbol{\xi}}\right)^{T} \cdot \boldsymbol{\eta}} = (\mathrm{Det}\,\mathbb{H}) e^{\left(\bar{\boldsymbol{\xi}}\right)^{T} \cdot \mathbb{H}^{-1} \cdot \boldsymbol{\xi}},$$

where  $\eta$ ,  $\bar{\eta}$ ,  $\xi$  and  $\bar{\xi}$  are independent Grassmann vectors.

(b) Use this result to show that

$$\langle \eta_{i1}\eta_{i2}\dots\eta_{im}\bar{\eta}_{jm}\dots\bar{\eta}_{j2}\bar{\eta}_{j1}\rangle = \sum_{\text{all pairings}} (-1)^P \mathbb{H}_{i1,jP_1}^{-1}\dots\mathbb{H}_{im,jP_m}^{-1},$$

where  $\{P_1, \ldots, P_m\}$  is a permutation of  $\{1, \ldots, m\}$ , P is the number of transpositions in this permutation, and the average is defined as in the previous problem.

- 3. Coherent states. Consider the coherent states  $a_i |\psi\rangle = \psi_i |\psi\rangle$ ,  $\langle \psi | a_i^{\dagger} = \langle \psi | \bar{\psi}_i$  with  $\psi$  and  $\bar{\psi}$  being independent (complex for bosons, and Grassmann for fermions) vectors. Show
  - (a) the completeness relation

$$\mathbb{I} = \int \mathrm{d} \left( \bar{\boldsymbol{\psi}}, \boldsymbol{\psi} \right) e^{-\bar{\boldsymbol{\psi}} \cdot \boldsymbol{\psi}} \ket{\boldsymbol{\psi}} \bra{\boldsymbol{\psi}},$$

(b) the trace

$$\operatorname{tr} \{A\} = \int \mathrm{d} \left( \bar{\psi}, \psi \right) e^{-\bar{\psi} \cdot \psi} \left\langle \zeta \psi \right| A \left| \psi \right\rangle$$

(c) and that, for a normal ordered operator  $O \equiv O(a_1^{\dagger}, \ldots, a_n^{\dagger}, a_1, \ldots, a_n)$ , the matrix element is

$$\langle \boldsymbol{\psi} | O | \boldsymbol{\psi}' \rangle = O(\bar{\psi}_1, \dots, \bar{\psi}_n, \psi_1', \dots, \psi_n') e^{\boldsymbol{\psi} \cdot \boldsymbol{\psi}'}$$

Here,  $\zeta = 1$  for bosons and  $\zeta = -1$  for fermions, and  $d(\bar{\psi}, \psi) = \prod_i d(\bar{\psi}_i, \psi_i)$ , where  $d(\bar{\psi}_i, \psi_i) = \pi^{-1} d(\Re \psi_i) d(\Im \psi_i)$  for bosons, and  $d(\bar{\psi}_i, \psi_i) = d\bar{\psi}_i d\psi_i$  for fermions.