## Oscillation with coercive force

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**Solution:** The particular solution can be found by the ansatz  $x_p(t) = be^{-t/\tau}$  with  $b = \beta \tau^2/(m - \gamma \tau + \kappa \tau^2)$ . Depending on whether  $\gamma^2 - 4\kappa m > = , < 0$  we have for a.  $\gamma^2 - 4\kappa m \equiv \Gamma^2 > 0$ 

$$x(t) = e^{-\frac{\gamma t}{2m}} \left( A e^{-\frac{\Gamma t}{2m}} + B e^{\frac{\Gamma t}{2m}} \right) + \frac{\beta \tau^2 m \ e^{-\frac{\tau}{\tau}}}{(m - \frac{\gamma \tau}{2})^2 - \frac{\Gamma^2 \tau^2}{4}} ;$$

for b.  $\gamma^2 - 4\kappa m = 0$ 

$$x(t) = e^{-\frac{\gamma t}{2m}} (A + Bt) + \frac{\beta \tau^2 m \ e^{-\frac{t}{\tau}}}{(m - \frac{\gamma \tau}{2})^2} ;$$

for c.  $\kappa m - \gamma^2 \equiv \Omega^2 > 0$ 

$$x(t) = e^{-\frac{\gamma t}{2m}} \left( A \cos \frac{\Omega t}{2m} + B \sin \frac{\Omega t}{2m} \right) + \frac{\beta \tau^2 m \ e^{-\frac{t}{\tau}}}{(m - \frac{\gamma \tau}{2})^2 + \frac{\Omega^2 \tau^2}{4}}$$

Through the boundary conditions x(0) = 0,  $\dot{x}(0) = 0$  we can determine A, B. Furthermore, we immediately see that  $x(\infty) = 0$ ,  $\dot{x}(\infty) = 0$  for all of the above mentioned solutions. Since the work of the two non-conservative forces equals the difference between the kinetic energy and the work for conservative forces, that is, A = 0, the entire work of F(t) has been dissipated by  $F_R$ .