Thermal population of a harmonic oscillator

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Solution: The density operator,

$$\hat{\rho} = \sum_{n} P_n |n\rangle \langle n| ,$$

satisfies $\hat{\rho}^2 \neq \hat{\rho}$ e Tr $\hat{\rho} = 1$. It allows to calculate the most likely value for the population and is obtained using the rule $\sum_{n=0}^{\infty} U^n = (1-U)^{-1}$ with the abbreviation $U \equiv e^{-\hbar \omega/k_B T}$,

$$\begin{split} \langle \hat{n} \rangle &= Tr \; \hat{\rho} \hat{n} = \sum_{m} \langle m | \hat{\rho} \hat{n} | m \rangle = \sum_{m} \langle m | \sum_{n} P_{n} | n \rangle \langle n | \hat{n} | m \rangle = \sum_{m} m P_{m} \langle m | m \rangle \\ &= \sum_{m} \frac{m e^{-m\beta\hbar\omega}}{\sum_{n} e^{-n\beta\hbar\omega}} = (1 - e^{-\beta\hbar\omega}) \sum_{m} m e^{-m\beta\hbar\omega} = (1 - e^{-\beta\hbar\omega}) \frac{-1}{\hbar\omega} \frac{\partial}{\partial\beta} \sum_{m} e^{-m\beta\hbar\omega} \\ &= (1 - e^{-\beta\hbar\omega}) \frac{-1}{\hbar\omega} \frac{\partial}{\partial\beta} \frac{1}{1 - e^{-\beta\hbar\omega}} = \frac{e^{-\beta\hbar\omega}}{1 - e^{-\beta\hbar\omega}} = \frac{1}{e^{\beta\hbar\omega} - 1} \; . \end{split}$$

For the average energy, with $E_n = n\hbar\omega$,

$$\langle \hat{E} \rangle = \sum_{n} E_{n} p_{n} = \frac{\hbar \omega}{e^{\beta \hbar \omega} - 1} .$$

This is precisely the distribution proposed by Planck for the light modes in the blackbody radiator. We can now express the occupation probability of state n as,

$$P_n = \frac{e^{-n\beta\hbar\omega}}{\sum_m e^{-m\beta\hbar\omega}} = (1 - e^{-\beta\hbar\omega})e^{-n\beta\hbar\omega} = \frac{\langle \hat{n} \rangle^n}{(1 + \langle \hat{n} \rangle)^{n+1}} .$$