

Bloch equations for three levels

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Solution: Let us first verify that the Bloch equations for two levels follow the formula (2.99). The Hamiltonian is $\hat{H} = \hbar\omega_0|2\rangle\langle 2| + \frac{\hbar}{2}\Omega e^{-i\omega t}|1\rangle\langle 2| + \frac{\hbar}{2}\Omega e^{i\omega t}|2\rangle\langle 1|$ the operator of Liouville is

$$\begin{aligned}\mathcal{L}\rho &= \sum_{i,j} \gamma_{ij} \left(|j\rangle\langle i| \sum_{m,n} |m\rangle\rho_{mn}\langle n|i\rangle\langle j| - \sum_{m,n} |m\rangle\rho_{mn}\langle n|i\rangle\langle j|j\rangle\langle i| \right. \\ &\quad \left. + |j\rangle\langle i| \sum_{m,n} |m\rangle\rho_{mn}\langle n|i\rangle\langle j| - |i\rangle\langle j|j\rangle\langle i| \sum_{m,n} |m\rangle\rho_{mn}\langle n| \right) \\ &= \gamma \left(|2\rangle\langle 1| \sum_{m,n} |m\rangle\rho_{mn}\langle n|1\rangle\langle 2| - \sum_{m,n} |n\rangle\rho_{nm}\langle m|1\rangle\langle 2|2\rangle\langle 1| \right. \\ &\quad \left. + |2\rangle\langle 1| \sum_{m,n} |m\rangle\rho_{mn}\langle n|1\rangle\langle 2| - |1\rangle\langle 2|2\rangle\langle 1| \sum_{m,n} |n\rangle\rho_{nm}\langle m| \right) \\ &= \gamma \left(2|2\rangle\rho_{11}\langle 2| - \sum_n |n\rangle\rho_{n1}\langle 1| - \sum_m |1\rangle\rho_{m1}\langle m| \right) .\end{aligned}$$

Follows the master equation,

$$\begin{aligned}\frac{d}{dt}\rho &= -\frac{i}{\hbar}[\hat{H}, \rho] + \mathcal{L}\rho \\ &= -i \left(\begin{array}{cc} \frac{1}{2}\Omega e^{-i\omega t}\rho_{21} - \frac{1}{2}\Omega e^{i\omega t}\rho_{12} & \frac{1}{2}\Omega e^{-i\omega t}\rho_{22} - \frac{1}{2}\Omega e^{-i\omega t}\rho_{11} - \omega_0\rho_{12} \\ \frac{1}{2}\Omega e^{i\omega t}\rho_{11} - \frac{1}{2}\Omega e^{i\omega t}\rho_{22} + \omega_0\rho_{21} & \frac{1}{2}\Omega e^{i\omega t}\rho_{12} - \frac{1}{2}\Omega e^{-i\omega t}\rho_{21} \end{array} \right) + \gamma \begin{pmatrix} 2\rho_{22} & -\rho_{12} \\ -\rho_{21} & -2\rho_{22} \end{pmatrix} \\ &= -i \left(\begin{array}{cc} \frac{\Omega}{2}\tilde{\rho}_{21} - \frac{\Omega}{2}\tilde{\rho}_{12} & \frac{\Omega}{2}e^{-i\omega t}\rho_{22} - \frac{\Omega}{2}\rho_{11}e^{-i\omega t} - \omega_0\tilde{\rho}_{12}e^{-i\omega t} \\ \frac{\Omega}{2}e^{i\omega t}\rho_{11} - \frac{\Omega}{2}e^{i\omega t}\rho_{22} + \omega_0\tilde{\rho}_{21}e^{i\omega t} & \frac{\Omega}{2}\tilde{\rho}_{12} - \frac{1}{2}\Omega\tilde{\rho}_{21} \end{array} \right) + \gamma \begin{pmatrix} 2\rho_{22} & -\rho_{12} \\ -\rho_{21} & -2\rho_{22} \end{pmatrix}\end{aligned}$$

transforming to the rotating frame, $\tilde{\rho}_{12}e^{-i\omega t} = \rho_{12}$,

$$\frac{d}{dt}\tilde{\rho} = \begin{pmatrix} -\frac{i}{2}\Omega\tilde{\rho}_{21} + \frac{i}{2}\Omega\tilde{\rho}_{12} & -\frac{i}{2}\Omega\rho_{22} + \frac{i}{2}\Omega\rho_{11} - i\Delta\tilde{\rho}_{12} \\ -\frac{i}{2}\Omega\rho_{11} + \frac{i}{2}\Omega\rho_{22} + i\Delta\tilde{\rho}_{21} & -\frac{i}{2}\Omega\tilde{\rho}_{12} + \frac{i}{2}\Omega\tilde{\rho}_{21} \end{pmatrix} + \gamma \begin{pmatrix} 2\rho_{22} & -\rho_{12} \\ -\rho_{21} & -2\rho_{22} \end{pmatrix} .$$

Moving to the Bloch vector notation we find the well-known expression (2.75).

We now generalize to three levels using the equations (2.100). For three levels the Hamiltonian is,

$$\hat{H} = \hbar\omega_2|2\rangle\langle 2| + \hbar\omega_3|3\rangle\langle 3| + \frac{\hbar}{2}\Omega_{12}e^{-i\omega_a t}|1\rangle\langle 2| + \frac{\hbar}{2}\Omega_{12}e^{i\omega_a t}|2\rangle\langle 1| + \frac{\hbar}{2}\Omega_{23}e^{i\omega_b t}|2\rangle\langle 3| + \frac{\hbar}{2}\Omega_{23}e^{-i\omega_b t}|3\rangle\langle 2| ,$$

and the Liouville operator is,

$$\begin{aligned}\mathcal{L}\rho &= \gamma_{12} (2|1\rangle\rho_{22}\langle 1| - 2|2\rangle\rho_{22}\langle 2| - |1\rangle\rho_{12}\langle 2| - |2\rangle\rho_{21}\langle 1|) \\ &\quad + \gamma_{23} (2|3\rangle\rho_{22}\langle 3| - 2|2\rangle\rho_{22}\langle 2| - |2\rangle\rho_{23}\langle 3| - |3\rangle\rho_{32}\langle 2|) \\ &\quad + \gamma_{13} (2|1\rangle\rho_{33}\langle 1| - 2|3\rangle\rho_{33}\langle 3| - |1\rangle\rho_{13}\langle 3| - |3\rangle\rho_{31}\langle 1|) .\end{aligned}$$

Follows the master equation,

$$\begin{aligned}
\frac{d}{dt}\rho &= -\frac{i}{\hbar}[\hat{H}, \rho] + \mathcal{L}\rho \\
&= -i \left[\begin{pmatrix} 0 & \frac{\Omega_{12}}{2}e^{-i\omega_a t} & 0 \\ \frac{\Omega_{12}}{2}e^{i\omega_a t} & -i\omega_2 & -\frac{i\Omega_{23}}{2}e^{i\omega_b t} \\ 0 & \frac{\Omega_{23}}{2}e^{-i\omega_b t} & -i\omega_3 \end{pmatrix}, \begin{pmatrix} \rho_{11} & \rho_{12} & \rho_{13} \\ \rho_{21} & \rho_{22} & \rho_{23} \\ \rho_{31} & \rho_{32} & \rho_{33} \end{pmatrix} \right] \\
&\quad + \begin{pmatrix} 2\gamma_{12}\rho_{22} + 2\gamma_{13}\rho_{33} & -\gamma_{12}\rho_{12} & -\gamma_{13}\rho_{13} \\ -\gamma_{12}\rho_{21} & -2(\gamma_{12} + \gamma_{23})\rho_{22} & -\gamma_{23}\rho_{23} \\ -\gamma_{13}\rho_{31} & -\gamma_{23}\rho_{32} & 2\gamma_{23}\rho_{22} - 2\gamma_{13}\rho_{33} \end{pmatrix}
\end{aligned}$$

Using MAPLE we derive the Bloch equations,

$$\begin{aligned}
\dot{\rho}_{11} &= \frac{i}{2}\Omega_{12} (e^{i\omega_a t} \rho_{12} - e^{-i\omega_a t} \rho_{21}) + 2\gamma_{12}\rho_{22} + 2\gamma_{13}\rho_{33} = 1 - \rho_{22} - \rho_{33} \\
\dot{\rho}_{22} &= -\frac{i}{2}\Omega_{12} (e^{i\omega_a t} \rho_{12} - e^{-i\omega_a t} \rho_{21}) + \frac{i}{2}\Omega_{23} (e^{-i\omega_b t} \rho_{23} - e^{i\omega_b t} \rho_{32}) - 2\gamma_{12}\rho_{22} - 2\gamma_{23}\rho_{22} \\
\dot{\rho}_{12} &= \frac{i}{2}\Omega_{12}e^{-i\omega_a t}(\rho_{11} - \rho_{22}) + \frac{i}{2}\Omega_{23}e^{-i\omega_b t}\rho_{13} + (i\omega_2 - \gamma_{12})\rho_{12} = \dot{\rho}_{21}^* \\
\dot{\rho}_{13} &= \frac{i}{2}\Omega_{23}e^{-i\omega_b t}\rho_{12} - \frac{i}{2}\Omega_{12}e^{-i\omega_a t}\rho_{23} + (i\omega_3 - \gamma_{13})\rho_{13} = \dot{\rho}_{31}^* \\
\dot{\rho}_{23} &= -\frac{i}{2}\Omega_{12}e^{i\omega_a t}\rho_{13} + \frac{i}{2}\Omega_{23}e^{i\omega_b t}(\rho_{22} - \rho_{33}) + (i\omega_3 - i\omega_2 - \gamma_{23})\rho_{23} = \dot{\rho}_{32}^* .
\end{aligned}$$

This result can be cast into a matrix form.